

**例 10.1**

一在基態上的氫原子，於時間  $t=0$  時，開始受到一個沿  $z$  軸方向的均勻電場作用，且其大小  $E = E_0 e^{-t/\tau}$ ，求當  $t \gg \tau$  時，此原子被激發至  $n=2$  狀態的機率為何？

**解** 氫原子與外加電場之交互作用位能為

$$V = -\vec{p} \cdot \vec{E} = qE\vec{r} \cdot \vec{e}_z = qE_0 e^{-t/\tau} Z$$

其對應之運算子則為

$$\hat{H}' = qE_0 e^{-t/\tau} \hat{Z} = \hat{H}'(t)$$

自基態  $n=1$  躍遷至  $n=2$  狀態之躍遷矩陣元為

$$\hat{H}'_{fi} = \langle \varphi_{n=2} | \hat{H}' | \varphi_{n=1} \rangle = qE_0 \langle \varphi_{n=2} | \hat{Z} | \varphi_{n=1} \rangle$$

其中  $\varphi_{n=2} = \varphi_{n\ell m} = \varphi_{200, 210, 21\pm 1}$ 、 $\varphi_{n=1} = \varphi_{n\ell m} = \varphi_{100}$ 。當末態  $n=2$  為  $\varphi_{200}$  時，由內積定義，及球面調和函數之正交性，故

$$\begin{aligned} \langle \varphi_{200} | \hat{Z} | \varphi_{100} \rangle &= \iint \langle \varphi_{200} | \vec{r}' \rangle \langle \vec{r}' | \hat{Z} | \vec{r}' \rangle \langle \vec{r}' | \varphi_{100} \rangle d^3 r' d^3 r \\ &= \iint \varphi_{200}^*(\vec{r}') Z \delta(\vec{r} - \vec{r}') \varphi_{100}(\vec{r}') d^3 r' d^3 r = \int \varphi_{200}^*(r) Z \varphi_{100}(r) d^3 r \\ &= \sqrt{\frac{4\pi}{3}} \int_0^\infty R_{20}(r) R_{10}(r) r^3 dr \int_{\Omega} Y_2^0 Y_1^0 Y_0^0 d\Omega = 0 \end{aligned}$$

同理，當末態為  $\varphi_{21\pm 1}$  時，因  $m_1 + m_2 + m_3 = \mp 1 + 0 + 0 = \mp 1 \neq 0$

$$\langle \varphi_{21\pm 1} | \hat{Z} | \varphi_{100} \rangle = 0$$

但當末態為  $\varphi_{210}$  時，

$$\langle \varphi_{210} | \hat{Z} | \varphi_{100} \rangle = \frac{1}{4\sqrt{2}\pi a^4} \int_0^\infty r e^{-r/2a} r e^{-r/a} r^2 dr \int_{\Omega} \cos\theta \cos\theta \sin\theta d\theta d\varphi$$

因  $\int_0^\infty r^4 e^{-3r/2a} dr = \frac{4!}{(3/2a)^5} = \left(\frac{4}{3}\right)^4 a^5$ ，故

$$\langle \varphi_{210} | \hat{Z} | \varphi_{100} \rangle = \frac{1}{4\sqrt{2}\pi a^4} = \left(\frac{4}{3}\right)^4 a^5 \cdot \frac{4\pi}{3} = \frac{1}{\sqrt{2}} \left(\frac{4}{3}\right)^4 a \cdot \frac{1}{3}$$

另外，含時間部分之積分項為

$$\begin{aligned} \int_0^t e^{i\omega_{21}t'} f(t') dt' &= \int_0^t e^{i\omega_{21}t'} e^{-t'/\tau} dt' = \int_0^t e^{\left(\frac{1}{\tau} + i\omega_{21}\right)t'} dt' = \int_0^t e^{ct'} dt' = \frac{1}{c} \int_0^t e^{ct'} d(ct') \\ &= \frac{1}{c} \int_0^{Z=ct} e^{Z'} dZ' = \frac{1}{c} [e^{Z'}]_0^{Z=ct} = \frac{1}{c} [e^{ct} - e^0] = \frac{1}{c} (e^{ct} - 1) \end{aligned}$$

其中  $c = -1/\tau + i\omega_{21}$ ，且

$$\omega_{21} = 1/\hbar(E_2 - E_1) = -13.6/\hbar(1/4 - 1) = 10.2/\hbar \quad (\text{eV})$$

因此

$$\begin{aligned} \left| \int_0^t e^{i\omega_{21}t'} e^{-t'/\tau} dt' \right|^2 &= \frac{1}{|c|^2} |e^{-t/\tau} e^{i\omega_{21}t} - 1|^2 = \frac{1}{|c|^2} |e^{-t/\tau} \cos \omega_{21}t + ie^{-t/\tau} \sin \omega_{21}t - 1|^2 \\ &= \frac{1}{(1/\tau)^2 + \omega_{21}^2} [(e^{-t/\tau} \cos \omega_{21}t - 1)^2 + (e^{-t/\tau} \sin \omega_{21}t)^2] \\ &= \frac{\tau^2 (e^{-2t/\tau} - 2e^{-t/\tau} \cos \omega_{21}t + 1)}{1 + \tau^2 \omega_{21}^2} \rightarrow \frac{\tau^2}{1 + \tau^2 \omega_{21}^2} \quad (\text{當 } t \gg \tau) \end{aligned}$$

最後，躍遷機率為

$$\begin{aligned} P_{if}(t) &= P_{100 \rightarrow 210}(t) \left| \frac{\hat{H}'_{fi}}{\hbar} \right|^2 \left| \int_0^t e^{i\omega_{21}t'} f(t') dt' \right|^2 \\ &= \frac{2^{15}}{3^{10}} \frac{q^2 E_0^2}{\hbar^2} a^2 \left| \int_0^t e^{i\omega_{21}t'} e^{-t'/\tau} dt' \right|^2 = \frac{2^{15}}{3^{10}} \frac{q^2 E_0^2 a^2}{\hbar^2} \frac{\tau^2}{1 + \tau^2 \omega_{21}^2} \end{aligned}$$

### 問題 10.5

在  $x=0$  至  $x=a$  之無限位能井中，有一粒子於  $t>0$  時，處於基態，當時間  $t=0$ ，開始受到微擾  $\hat{V}(t)=\varepsilon e^{-t^2}$  之作用，其中  $\varepsilon \ll 1$ 。計算在時間  $t=+\infty$  時，粒子被發現處於第一激發態的機率是多少？

**解** 由基態  $n=1$  到第一激發態的躍遷機率為

$$P_{1 \rightarrow 2} = \frac{1}{\hbar^2} \left| \int_0^\infty \langle \psi_2 | \hat{V}(t) | \psi_1 \rangle e^{i\omega_{21}t} dt \right|^2$$

其中 
$$\omega_{21} = \frac{E_2 - E_1}{\hbar} = \frac{4\pi^2 \hbar}{2ma^2} - \frac{\pi^2 \hbar}{2ma^2} = \frac{3\pi^2 \hbar}{2ma^2}$$

$$\langle \psi_2 | \hat{V}(t) | \psi_1 \rangle = \frac{2\varepsilon}{a} e^{-t^2} \int_0^a \sin\left(\frac{2\pi x}{a}\right) \sin\left(\frac{\pi x}{a}\right) dx = \frac{16\varepsilon a}{9\pi^2} e^{-t^2}$$

此處我們用到  $E_n = \frac{n^2 \pi^2 \hbar^2}{2ma^2}$  以及  $\psi_n = \sqrt{\frac{2}{a}} \sin\left(\frac{n\pi x}{a}\right)$

故 
$$P_{1 \rightarrow 2} = \left( \frac{16\varepsilon a}{9\pi^2 \hbar^2} \right)^2 \left| \int_0^\infty e^{i\omega_{21}t - t^2} dt \right|^2$$

利用變數變換  $y = t - \frac{i}{2}\omega_{21}$ ，使得  $i\omega_{21}t - t^2 = -\frac{\omega_{21}^2}{4} - y^2$  且  $dt = dy$ 。

$$P_{1 \rightarrow 2} = \left( \frac{16\varepsilon a}{9\pi^2 \hbar^2} \right)^2 \left| e^{-\omega_{21}^2/4} \int_0^\infty e^{-y^2} dy \right|^2 = \left( \frac{8\varepsilon a}{9\pi \hbar} \right)^2 \frac{1}{\pi} e^{-\omega_{21}^2/2}$$

上式使用了  $\int_0^\infty e^{-y^2} dy = \frac{\sqrt{\pi}}{2}$  以及  $\omega_{21} = \frac{3\pi^2 \hbar^2}{2ma^2}$